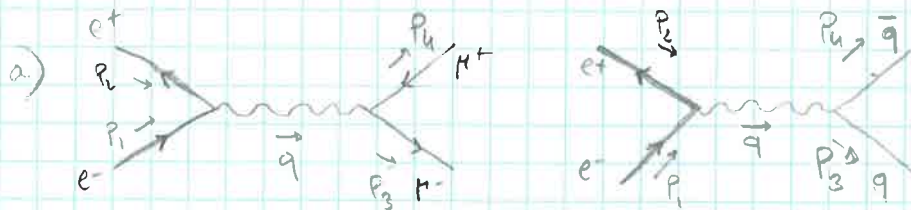


# EXERCISE SHEET 4 - KT 11

## 1) ELECTRON-POSITRON ANNIHILATION

$$e^+e^- \rightarrow \mu^+\mu^- \quad , \quad e^+e^- \rightarrow q\bar{q}$$



b)  $e^+e^- \rightarrow \mu^+\mu^-$

$$-iM = (2\pi)^4 \int d^4q \delta^4(p_1 + p_2 - q) \delta^4(p_3 + p_4 - q) \\ \times \bar{v}(2) i g_e \gamma^\mu u(1) \frac{i g_\mu^{\nu}}{q^2} \bar{u}(3) i g_e \gamma^\nu v(4)$$

$$-iM = \frac{g_e^2}{(p_1 + p_2)} [\bar{v}(2) \gamma^\mu u(1)] [\bar{u}(3) \gamma_\mu v(4)]$$

$e^+e^- \rightarrow q\bar{q}$  → the matrix element is the same as for  $e^+e^- \rightarrow \mu^+\mu^-$  but considering that the charge carried out by quarks is only a fraction of that carried by an electron, so

$$-iM = \frac{Q_q g_e^2}{p_1 + p_2} [\bar{v}(2) \gamma^\mu u(1)] [\bar{u}(3) \gamma_\mu v(4)] \times 3$$

where  $Q_q$  indicates the charge of the quark.

So  $M(e^+e^- \rightarrow q\bar{q}) = Q_q M(e^+e^- \rightarrow \mu^+\mu^-)$

c) the total cross section of  $e^+e^- \rightarrow q\bar{q}$  is equal to

$$\sigma(e^+e^- \rightarrow q\bar{q}) = 3 Q_q^2 \sigma(e^+e^- \rightarrow \mu^+\mu^-)$$

where the factor 3 takes into account the fact that we have a diagram per each colour of the quark.

Summing up over the different flavour quarks we get

$$\sigma(e^+e^- \rightarrow \text{hadrons}) = \sum_q \sigma(e^+e^- \rightarrow q\bar{q}) = 3 \sum_q Q_q^2 \sigma(e^+e^- \rightarrow \mu^+\mu^-)$$

$$\text{where } \sigma(e^+e^- \rightarrow \mu^+\mu^-) = \frac{4\pi\alpha^2}{3s^2}$$

$$R = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = 3 \sum_q Q_q^2$$

d)  $R_S = R(1 + \alpha_s/\pi)$

At  $E_{cm} = 2.8 \text{ GeV} \rightarrow u, d, s$  quarks can be produced

$E_{cm} = 5 \text{ GeV} \rightarrow u, d, s, c$

$E_{cm} = 15 \text{ GeV} \rightarrow u, d, s, c, b$

For  $E_{cm} = 2.8 \text{ GeV}$   $R = 3 \left[ \left(\frac{2}{3}\right)^2 + \left(\frac{1}{3}\right)^2 + \left(\frac{1}{3}\right)^2 \right] = 2$  for  $u, d, s$

For  $E_{cm} = 5 \text{ GeV}$   $R = 2 + 3 \left(\frac{2}{3}\right)^2 = \frac{10}{3}$  for  $u, d, s, c$

For  $E_{cm} = 15 \text{ GeV}$   $R = \frac{10}{3} + 3 \left(\frac{1}{3}\right)^2 = \frac{11}{3}$  for  $u, d, s, c, b$

$\alpha_s(\mu^2 = E_{cm}^2) = 0.27 \times E_{cm} = 2.8 \text{ GeV}$

$0.2 \times E_{cm} = 5 \text{ GeV}$

$0.18 \times E_{cm} = 15 \text{ GeV}$

$R_S = R(1 + \alpha_s/\pi)$

$R_S(E_{cm} = 2.8 \text{ GeV}) \approx 2.17$

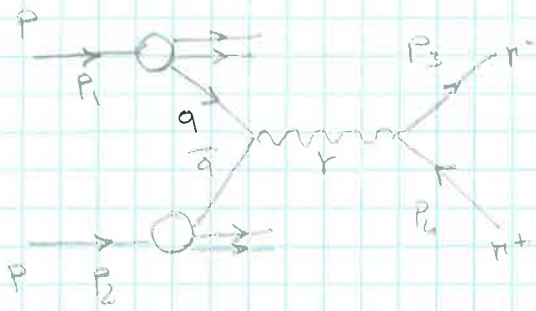
$R_S(E_{cm} = 5 \text{ GeV}) \approx 3.54$

$R_S(E_{cm} = 15 \text{ GeV}) \approx 3.88$

$R = \frac{11}{3} + 3 \left(\frac{2}{3}\right)^2 = \frac{15}{3} = 5$



# DRELL YAN PROCESS



$$PP \rightarrow \mu^+ \mu^- X$$

$$\sigma(e^+e^- \rightarrow \mu^+\mu^-) = \frac{4\pi\alpha^2}{3s}$$

$$\sigma(q\bar{q} \rightarrow \mu^+\mu^-) = \frac{1}{N_c} Q_q^2 \frac{4\pi\alpha^2}{3s} \rightarrow \frac{1}{N_c} = \frac{1}{3}$$

This factor accounts for the fact that of the 9 possible combinations of  $q\bar{q}$ , the annihilation process can take place only for 3  $\rightarrow \mu^+\mu^-$ ,  $q\bar{q}$  and  $\bar{q}q$ .

Quark and anti-quark bring only a fraction of the proton, namely  $x_1$  and  $x_2 \Rightarrow x_1 + dx_1$  and  $x_2 + dx_2$

$$d^2\sigma = Q_q^2 \frac{4\pi\alpha^2}{3s} u^p(x_1) dx_1 \bar{u}^p(x_2) dx_2$$

$\hookrightarrow$  this comes from the definition of parton distribution function, PDF

$\bar{u}^p$  is the PDF for the anti-up quark in the proton. The PDF of quarks in the proton are not the same as PDF of anti-quarks in proton.

The centre of mass energy of the  $q\bar{q}$  system can be written in terms of initial momenta of the protons:

$$\hat{s} = (x_1 P_1 + x_2 P_2)^2 = x_1^2 P_1^2 + x_2^2 P_2^2 + 2x_1 x_2 P_1 \cdot P_2$$

In high energy limit  $P_1^2 = P_2^2 = M_p^2 \approx 0 \Rightarrow \hat{s} \approx 2x_1 x_2 P_1 \cdot P_2$

$\Xi \approx x_1, x_2$  ( $2 P_1, P_2$ ) =  $x_1, x_2 S \rightarrow$  S is the centre of mass of the colliding pp system.

$$d^2G = \frac{4}{9} \cdot \frac{4\pi d^2}{9x_1 x_2 S} u^P(x_1) \bar{u}^P(x_2) dx_1 dx_2$$

We also have a contribution from  $d\bar{d}$  annihilation

$$d^2G = \frac{4\pi d^2}{9x_1 x_2 S} \left[ \frac{4}{9} \{ u^P(x_1) \bar{u}^P(x_2) \} + \frac{1}{9} \{ d^P(x_1) \bar{d}^P(x_2) \} + \frac{4}{9} \{ \bar{u}^P(x_1) u^P(x_2) \} + \frac{1}{9} \{ \bar{d}^P(x_1) d^P(x_2) \} \right] dx_1 dx_2$$

$\rightarrow$  We do not have just valence quarks but other gluon exchange that can fluctuate into virtual  $q\bar{q}$  pairs. These sea and quark gluons are suppressed by a factor  $1/q^2$ , so they are produced mainly at low  $x$ .

$\Rightarrow$  This means we have  $u(x) = u_v(x) + u_s(x) \rightarrow$  accounting for the contribution from valence and sea quarks

Concerning antiquarks  $\rightarrow$  then we only see quark contributions

$$\bar{u}(x) = \bar{u}_s(x) \quad \text{and} \quad \bar{d}(x) = \bar{d}_s(x)$$

Given that the proton consists of 2 valence quarks and one valence down quark the valence quark PDFs are normalized differently

$$\int_0^1 u_v(x) dx = 2 \quad \int_0^1 d_v(x) dx = 1$$

Few assumptions  $q_s(x) = \bar{q}_s(x)$  because sea quarks and antiquarks are produced in couple

Moreover  $u_u \sim u_d = 0$  so PDFs can be assumed to be approximately the same

$$u_s(x) = \bar{u}_s(x) = d_s(x) = \bar{d}_s(x) \approx S(x)$$



Due to the proton the normalization factor for valence up quarks is 2

$$\Rightarrow D u_v(x) = 2 d_v(x)$$

But, let's write the cross section in terms of valence and sea quarks.

$$\begin{aligned} d^2 \sigma_{pp}^{\text{PP}} &= \frac{4\pi \alpha^2}{9x_1 x_2 S} \frac{1}{9} \left\{ 4 u_v(x_1) \bar{u}_s(x_2) + d_v(x_1) \bar{d}_s(x_2) + 4 \bar{u}_s(x_1) u_v(x_2) \right. \\ &\quad \left. + \bar{d}_s(x_1) d_v(x_2) + \text{the same with } u_s(x) \text{ and } d_s(x) \right\} \\ &= \frac{4\pi \alpha^2}{81 x_1 x_2 S} \left\{ 8 d_v(x_1) S(x_2) + d_v(x_1) S(x_2) + 8 S(x_1) d_v(x_2) \right. \\ &\quad \left. + S(x_1) d_v(x_2) + 10 S(x_1) S(x_2) \right\} \\ &= \frac{4\pi \alpha^2}{81 x_1 x_2 S} \left\{ 9 d_v(x_1) S(x_2) + 9 S(x_1) d_v(x_2) + 10 S(x_1) S(x_2) \right\} \end{aligned}$$

Due to the case of  $p\bar{p}$  collisions both  $\bar{u}$  and  $\bar{d}$  can be of valence, and we have

$$\begin{aligned} d^2 \sigma_{p\bar{p}}^{\text{PP}} &= \frac{4\pi \alpha^2}{81 x_1 x_2 S} \left\{ 17 d_v(x_1) d_v(x_2) + 9 d_v(x_1) S(x_2) + 9 S(x_1) d_v(x_2) + \right. \\ &\quad \left. + 10 S(x_1) S(x_2) \right\} dx_1 dx_2 \end{aligned}$$

$$\begin{aligned} \text{c)} \quad k &= \frac{\sigma(\bar{u}^+ C \rightarrow \pi^+ \pi^- X)}{\sigma(\bar{u}^- C \rightarrow \pi^+ \pi^- X)} & k &\rightarrow 1 \text{ for small } Q^2/S \\ & & k &\rightarrow 1/4 \text{ for } Q^2/S \rightarrow 1 \end{aligned}$$

$$\bar{u}^+ \rightarrow u \bar{d} \quad / \quad \bar{u}^- \rightarrow \bar{u} d$$

the PDFs for the  $e$  pairs can be written in terms of valence and sea quarks:

$$u^{\pi^+}(x) = u_v^{\pi^+}(x) + S^{\pi^+}(x) = u_v^{\pi^+}(x) + S^{\pi^+}(x)$$

$$\bar{d}^{\pi^+}(x) = \bar{d}_v^{\pi^+}(x) + S^{\pi^+}(x) = \bar{d}_v^{\pi^+}(x) + S^{\pi^+}(x)$$

$$\bar{u}^{\pi^+}(x) = S^{\pi^+}(x) = S^{\pi^+}(x)$$

$$d^{\pi^+}(x) = S^{\pi^+}(x) = S^{\pi^+}(x)$$

Assuming isospin symmetry

the down-quark PDF in  $\pi^-$

is identical to the up-quark

PDF in  $\pi^+$

$$u^{\pi^-}(x) = S^{\pi^-}(x)$$

$$\bar{d}^{\pi^-}(x) = S^{\pi^-}(x)$$

$$\bar{u}^{\pi^-}(x) = \bar{u}_V^{\pi^-}(x) + S^{\pi^-}(x)$$

$$d^{\pi^+}(x) = d_V^{\pi^+}(x) + S^{\pi^+}(x)$$

When  $Q^2/s \rightarrow 0 \Rightarrow x_1 \approx x_2 \approx 0 \Rightarrow$  sea quarks dominate and there is no difference between  $\pi^+C$  and  $\pi^-C$ .

When  $Q^2/s \approx 1$   $x_1 \approx x_2 \approx 1$  and the valence quarks dominate.

Put assuming isospin symmetry  $[d^{\pi^-}(x) = u^{\pi^+}(x)]$  and considering that  $C$  has the same number of up and down quarks  $\Rightarrow$

$$\frac{\sigma(\pi^+C \rightarrow p^+p^-x)}{\sigma(\pi^-C \rightarrow p^+p^-x)} = \frac{e_d^2 f_d^{\pi^+} f_d^C}{e_u^2 f_u^{\pi^-} f_u^C} = \frac{e_d^2}{e_u^2} = \frac{1}{4}$$

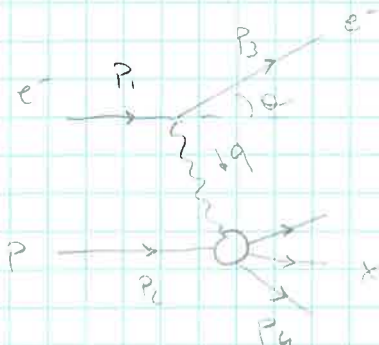
$\Rightarrow$   $d$  and  $u$  are valence quarks in  $C$ ,  
but  $\bar{u}$  and  $\bar{d}$  are sea quarks.



# ELECTRON-PROTON INELASTIC SCATTERING

$$X = \frac{Q^2}{2P_2 \cdot q} \quad \nu = \frac{P_2 \cdot q}{M_p}$$

$$\Rightarrow 2M_p \nu = W^2 + Q^2 - M_p^2$$



$Q^2 = -q^2 \rightarrow$  negative four-momentum squared of the virtual photon

$W$  is the four-momentum of the hadronic system ( $X$ )

$$W^2 \equiv P_4^2 = (P_1 + P_2)^2 = q^2 + 2P_2 \cdot q + P_2^2 = -Q^2 + 2P_2 \cdot q + M_p^2$$

$$\text{but } \nu = \frac{P_2 \cdot q}{M_p} \Rightarrow W^2 + Q^2 - M_p^2 = 2\nu M_p$$

In the rest frame of the proton  $P_2 = (M_p, 0, 0, 0)$ , the momenta of the initial state  $e^-$ , the final state  $e^-$  and the virtual photon can be written as:

$$P_1 = (E_1, 0, 0, E_1), \quad P_3 = (E_3, E_3 \sin\theta, 0, E_3 \cos\theta)$$

$$\text{and } q = (E_1 - E_3, \vec{P}_1 - \vec{P}_3)$$

$$\nu = \frac{P_2 \cdot q}{M_p} = \frac{1}{M_p} M_p (E_1 - E_3) = E_1 - E_3 \Rightarrow \text{this is simply the energy lost by the electron.}$$