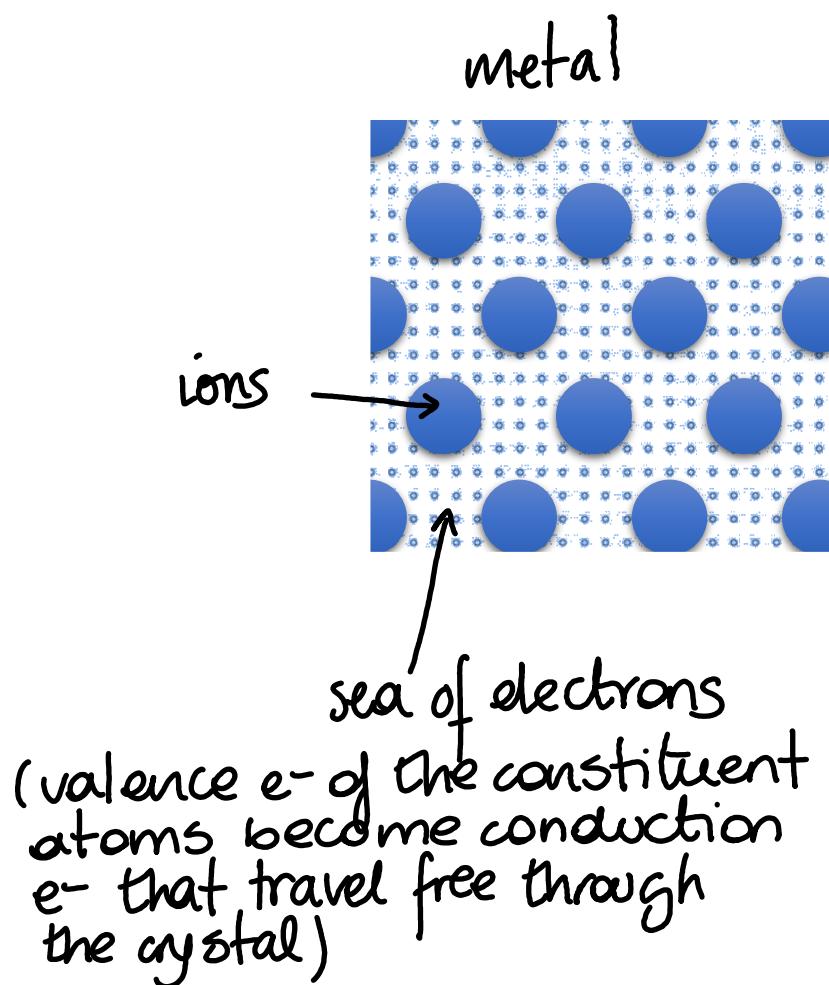


Free electron model

Kittel ch. 6

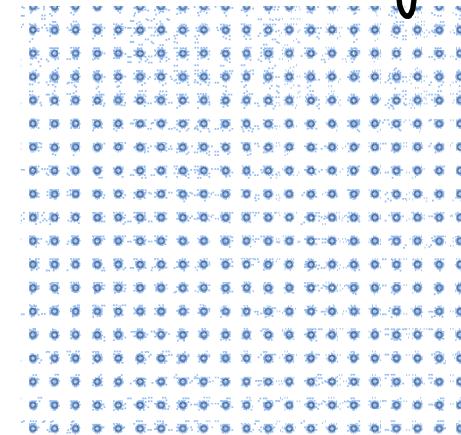
Starting Point for Understanding Electrons in Solids

Metals



application of Kinetic theory of gases to metals by considering them as a "gas of electrons"
[but e- must obey Pauli exclusion principle!]

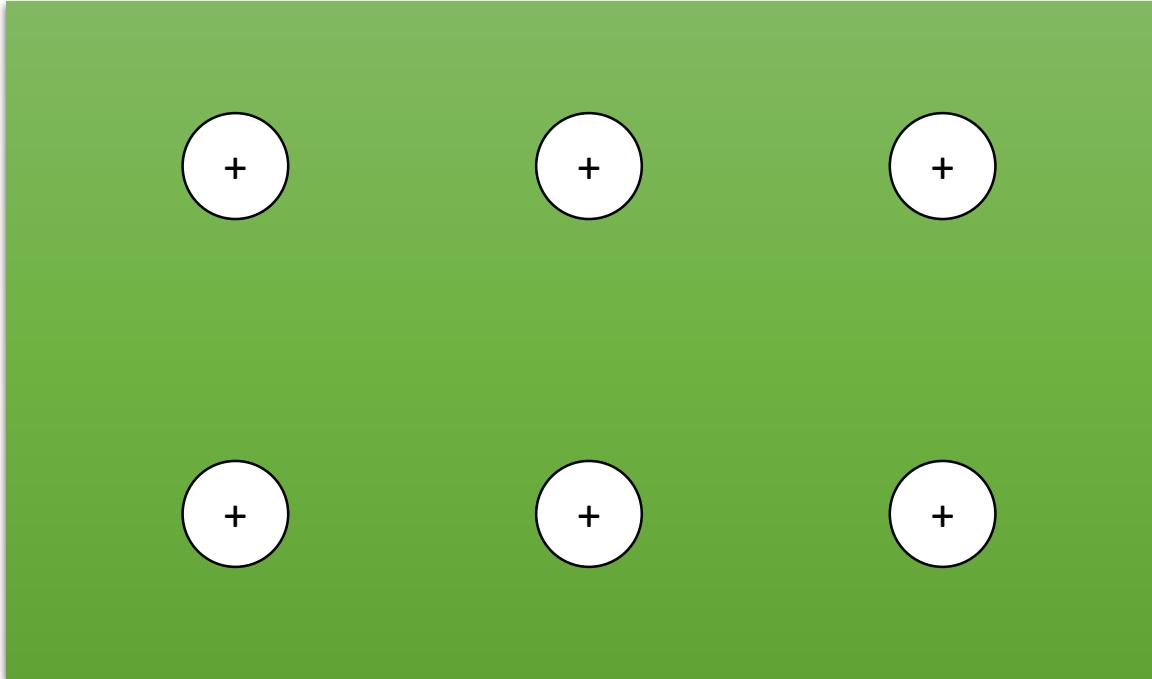
Free electron gas



- e- are free to move through crystal
- nuclei are replaced by smooth background

Alkali metals – How many conduction electrons?

3	Li	Lithium
	6.941	
11	Na	Sodium
	22.990	
19	K	Potassium
	39.098	
37	Rb	Rubidium
	84.468	



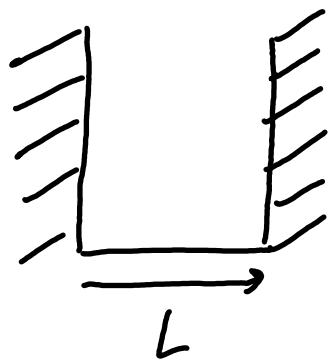
Na

Free atom: $1s^2 2s^2 2p^6 3s^1$

In the metal, this e^- becomes a conduction e^-

Monovalent crystal : N atoms \Rightarrow N conduction electrons

Free electrons in 1D



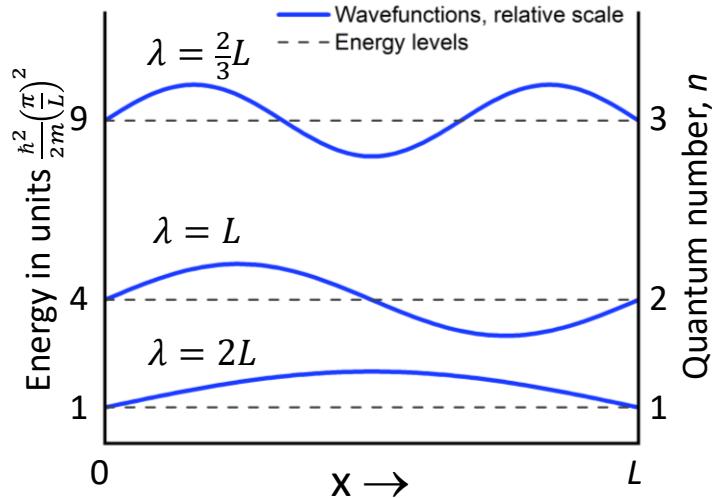
1 e⁻ of mass m confined to a length L by ∞ -potential barriers

$$\mathcal{H} \Psi_n(x) = E_n \Psi_n(x)$$

$$-\frac{\hbar^2}{2m} \frac{d^2\Psi_n}{dx^2} = \underbrace{E_n \Psi_n(x)}_{\text{energy of the e- in the orbital}}$$

boundary conditions $\Psi(0) = \Psi(L) = 0$

solution $\Psi_n(x) = A \cdot \sin\left(\frac{n\pi}{L} \cdot x\right)$ $n = \text{integer}$
 ↑
 constant



$$\text{Eigenvalues (Energy)} \quad E = \frac{\hbar^2}{2m} \left(\frac{n\pi}{L}\right)^2$$

Free electrons in 1D - electron occupancy

N valence electrons

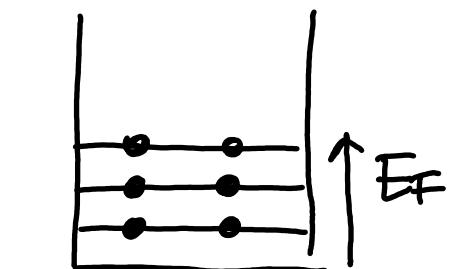
- must obey Pauli exclusion principle
- electronic state in 1D: quantum numbers n (describes $\psi_n(x)$)
 $m_s = \pm \frac{1}{2}$

- each orbital label n accommodates $2e^-$

- n_F : highest filled energy level

$$2 \cdot n_F = N$$

→ Fermi Energy = Energy of the highest occupied orbital



$$E_F = \frac{\hbar^2}{2m} \left(\frac{\pi N}{2L} \right)^2$$

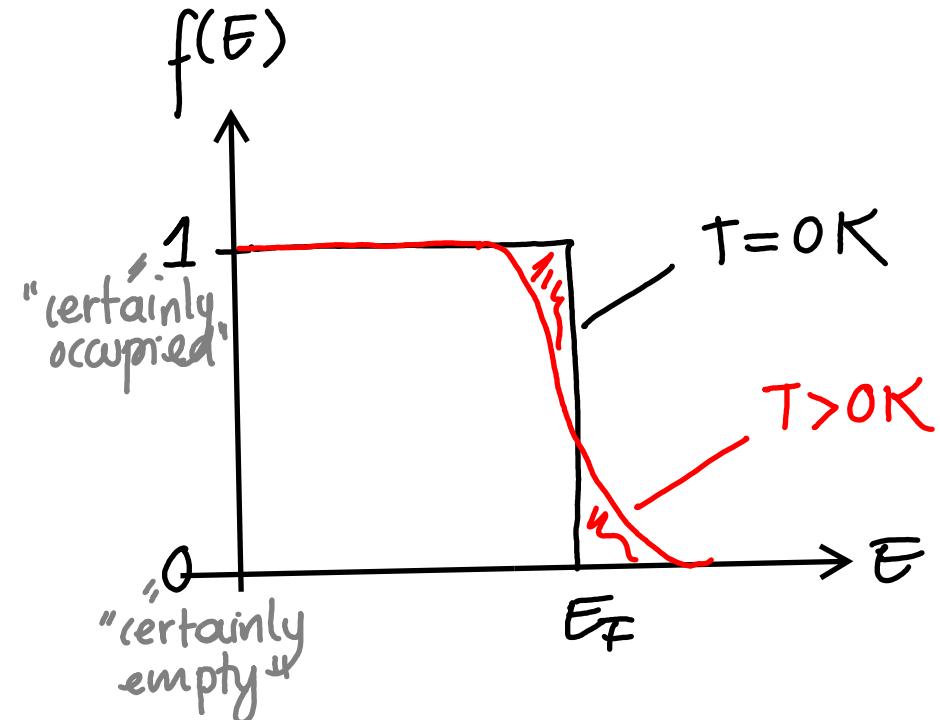
→ in metals $E_F \approx 5\text{ eV}$

Fermi distribution

$f(E)$ = distribution function
= distribution of e^- among levels

As $T \uparrow$, Kinetic energy of $e^- \uparrow$
 \Rightarrow only e^- close to E_F can be excited (Pauli Principle!)

- Energy an e^- absorb thermally $\sim k_B T$
(~ 0.025 eV at RT)
- $E_F \sim 5$ eV

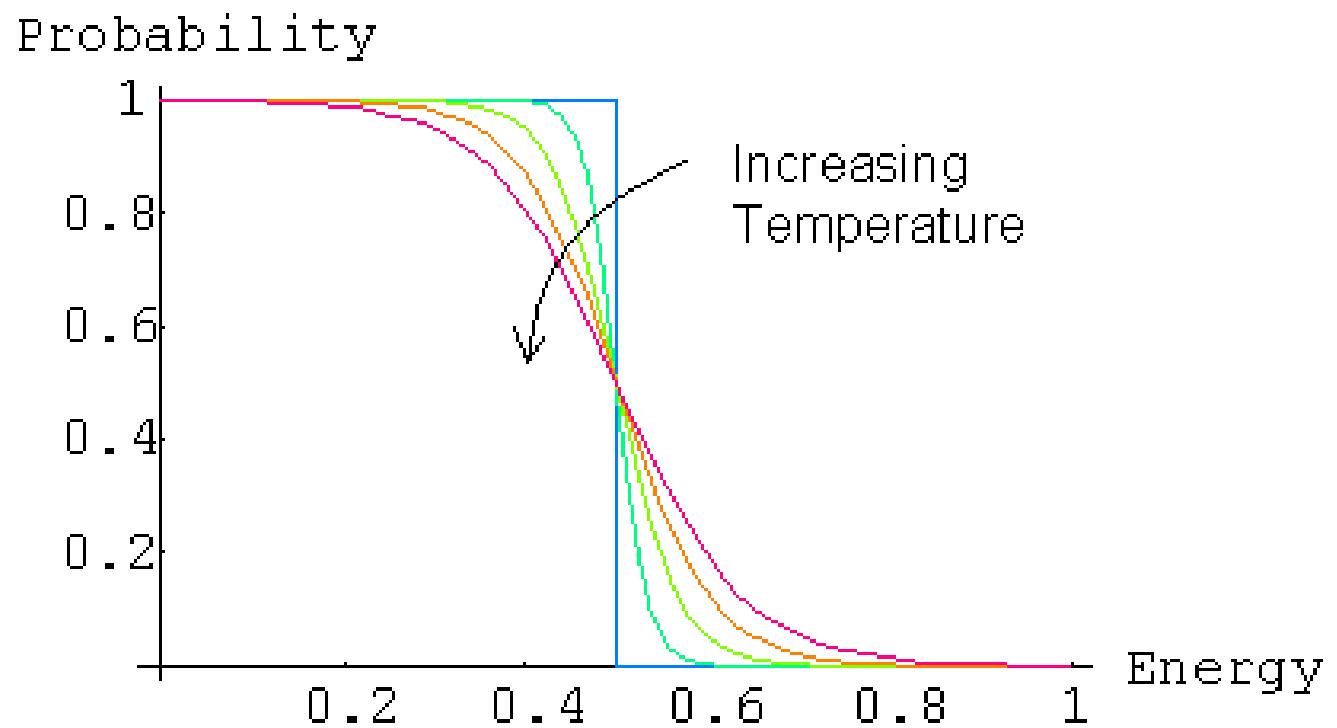


Fermi-Dirac distribution

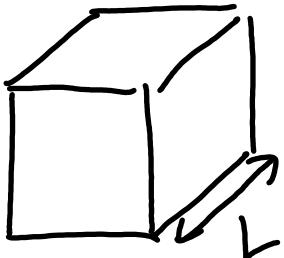
probability that an orbital of Energy E is occupied at thermal equilibrium.

$$f(E) = \frac{1}{\exp\left(\frac{E-\mu}{k_B T}\right) + 1}$$

μ : chemical potential
(at 0K: $\mu = E_F$)



Free electrons in 3D



$$\mathcal{H}\psi(\vec{r}) = -\frac{\hbar^2}{2m} \nabla^2 \psi(\vec{r}) = -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \psi(\vec{r}) = E \psi(\vec{r})$$

- periodic boundary conditions i.e. x -axis $\psi(x+L, y, z) = \psi(x, y, z)$
(idem for other axes)

- Solutions: $\psi_{\vec{k}}(\vec{r}) = A \exp(i\vec{k} \cdot \vec{r})$

with $k_i = \frac{2\pi}{L} n_i$ (i.e. $k_x = 0, \pm \frac{2\pi}{L}, \pm \frac{4\pi}{L}, \dots$)

n_i : positive integers

- Energies for the orbital
with K

$$E_K = \frac{\hbar^2 K^2}{2m} = \frac{\hbar^2}{2m} (k_x^2 + k_y^2 + k_z^2)$$

Free electrons in 3D

$$\psi_{\vec{k}} = A \exp(i \vec{k} \cdot \vec{r})$$

* $\psi(\vec{r})$ are eigenfunctions of the momentum

$$\vec{p} \psi_{\vec{k}}(\vec{r}) = -i\hbar \nabla \psi_{\vec{k}}(\vec{r}) = \hbar \vec{k} \psi_{\vec{k}}(\vec{r})$$

eigenvalue of the momentum

* velocity of the electron

$$\vec{v} = \frac{\vec{p}}{m} = \frac{\hbar \vec{k}}{m}$$

- Each \vec{k} state occupies $(\frac{2\pi}{L})^3$ in 3D

Here 2D example

